## Molar Specific Heats of an Ideal Gas

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## Molar Specific Heats

- The energy required to raise the temperature of $n$ moles of gas from $T_{i}$ to $T_{f}$ depends on the path taken between the initial and final states.
- Below is a PV diagram showing an isovolumetric process and an isobaric process that connect the same two isotherms.
- The energy transfer for the two processes can be written as

$$
\begin{aligned}
& Q=n C_{V} \Delta T \quad(\text { constant volume }) \\
& Q=n C_{p} \Delta T \quad(\text { constant pressure })
\end{aligned}
$$

where $C_{V}$ is the molar specific heat at constant volume and $C_{P}$ is the molar specific heat at constant pressure.


## Isovolumetric Processes

- The first law of thermodynamics can be written in differential form as

$$
d E_{i n t}=d Q+d W=d Q-P d V
$$

- In an isovolumetric process $d V=0$, so we have

$$
d E_{\text {int }}=d Q=n C_{V} d T
$$

- Thus, the change in internal energy be can be expressed as

$$
\Delta E_{i n t}=n C_{V} \Delta T
$$

- Since the change in internal energy is a state variable it does not depend on the path taken between the initial and final states, and the above equation holds for any process in which the temperature change is $\Delta T$.
- Also, the above equation is valid for monatomic, diatomic, and polyatomic gases.
- The differential form of the equation can be used to express the molar specific heat at constant volume as

$$
C_{V}=\frac{1}{n} \frac{d E_{i n t}}{d T}
$$

## Isobaric Processes

- Applying the first law to the isobaric process yields

$$
d E_{i n t}=d Q-P d V=n C_{P} d T-P d V
$$

- Using the ideal gas law, we can write $P d V=n R d T$
- Substituting this and $d E_{\text {int }}=n C_{V} d T$ into the expression for the first law gives

$$
\begin{gathered}
n C_{V} d T=n C_{p} d T-n R d T \\
C_{P}-C_{V}=R
\end{gathered}
$$

- This expression applies to any ideal gas and shows that the molar specific heat at constant pressure is greater than the molar specific heat at constant volume by the amount $R$


## Adiabatic Processes

- In an adiabatic process $d Q=0$, so the first laws gives us

$$
\begin{gathered}
d E_{\text {int }}=-P d V \\
d E_{\text {int }}+P d V=0 \\
n C_{V} d T+\frac{n R T}{V} d V=0 \\
d T+\frac{R}{C_{V}} T \frac{d V}{V}=0 \\
d T+\frac{C_{P}-C_{V}}{C_{V}} T \frac{d V}{V}=0 \\
d T+\left(\frac{C_{P}}{C_{V}}-1\right) T \frac{d V}{V}=0 \\
d T+(\gamma-1) T \frac{d V}{V}=0
\end{gathered}
$$

where $\gamma=\frac{C_{P}}{C_{V}}$.

- Multiplying by $V^{\gamma-1}$ we have

$$
\begin{gathered}
V^{\gamma-1} d T+T(\gamma-1) V^{\gamma-2} d V=0 \\
d\left(T V^{\gamma-1}\right)=0 \\
T V^{\gamma-1}=\text { constant }
\end{gathered}
$$

## Adiabatic Process II

- From the previous expression, we can write

$$
\begin{aligned}
T_{i} V_{i}^{\gamma-1} & =T_{f} V_{f}^{\gamma-1} \\
\frac{P_{i} V_{i}}{n R} V_{i}^{\gamma-1} & =\frac{P_{f} V_{f}}{n R} V_{f}^{\gamma-1} \\
P_{i} V_{i}^{\gamma} & =P_{f} V_{f}^{\gamma} \\
P V^{\gamma} & =\text { constant }
\end{aligned}
$$

- The PV diagram for an adiabatic expansion is shown below

Serway/Jewett; Principles of Physics, 3/e Figure 17.14


## Equipartition of Energy

- Consider a monatomic ideal gas undergoing an isovolumetric process
- We have shown that for an isovolumetric process the first law of thermodynamics becomes

$$
Q=\Delta E_{i n t}
$$

and that for a monatomic gas

$$
E_{\text {int }}=\frac{3}{2} n R T
$$

- So, the first law yields

$$
\begin{gathered}
n C_{V} \Delta T=\frac{3}{2} n R \Delta T \\
C_{V}=\frac{3}{2} R=12.5 \mathrm{~J} / \mathrm{mol} \cdot K
\end{gathered}
$$

- This prediction is in excellent agreement with measured values of $C_{V}$ for monatomic gases, but fails for diatomic and polyatomic gases

Serway/Jewett; Principles of Physics, 3/e
Table 17.3

| TAB | 17. | Molar Specific Heats of Various Gases |  |  |
| :---: | :---: | :---: | :---: | :---: |
|  | Molar Specific Heat $(\mathrm{J} / \mathrm{mol} \cdot \mathrm{K})^{\text {a }}$ |  |  |  |
|  | $C_{P}$ | $C_{V}$ | $C_{P}-\mathrm{C}_{V}$ | $\gamma=C_{P} / C_{V}$ |
| Monatomic Gases |  |  |  |  |
| He | 20.8 | 12.5 | 8.33 | 1.67 |
| Ar | 20.8 | 12.5 | 8.33 | 1.67 |
| Ne | 20.8 | 12.7 | 8.12 | 1.64 |
| Kr | 20.8 | 12.3 | 8.49 | 1.69 |
| Diatomic Gases |  |  |  |  |
| $\mathrm{H}_{2}$ | 28.8 | 20.4 | 8.33 | 1.41 |
| $\mathrm{N}_{2}$ | 29.1 | 20.8 | 8.33 | 1.40 |
| $\mathrm{O}_{2}$ | 29.4 | 21.1 | 8.33 | 1.40 |
| CO | 29.3 | 21.0 | 8.33 | 1.40 |
| $\mathrm{Cl}_{2}$ | 34.7 | 25.7 | 8.96 | 1.35 |
| Polyatomic Gases |  |  |  |  |
| $\mathrm{CO}_{2}$ | 37.0 | 28.5 | 8.50 | 1.30 |
| $\mathrm{SO}_{2}$ | 40.4 | 31.4 | 9.00 | 1.29 |
| $\mathrm{H}_{2} \mathrm{O}$ | 35.4 | 27.0 | 8.37 | 1.30 |
| $\mathrm{CH}_{4}$ | 35.5 | 27.1 | 8.41 | 1.31 |

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## Equipartition of Energy II

- For diatomic and polyatomic molecules, we must take into account the other types of motion that exists in addition to translation
- The figure below shows the possible motions of a diatomic molecule: (a) translational, (b) rotational, and (c) vibrational

Serway/Jewett; Principles of Physics, 3/e Figure 17.15


## Equipartition of Energy III

- The theorem of equipartition of energy states that each degree of freedom, on average, contributes $\frac{1}{2} k_{B} T$ of energy per molecule
- If we consider the rotational motion of a diatomic molecule in addition to its translational motion, we get 5 degrees of freedom (rotation about the y axis shown below can be ignored because the moment of inertia is so small) which for a gas with $N$ molecules yields

$$
\begin{gathered}
E_{\text {int }}=5 N\left(\frac{1}{2} k_{B} T\right)=\frac{5}{2} N k_{B} T=\frac{5}{2} n R T \\
C_{V}=\frac{1}{n} \frac{d E_{\text {int }}}{d T}=\frac{5}{2} R=20.8 \mathrm{~J} / \mathrm{mol} \cdot K \\
C_{P}=C_{V}+R=\frac{7}{2} R \\
\gamma=\frac{C_{p}}{C_{V}}=\frac{\frac{7}{2} R}{\frac{5}{2} R}=\frac{7}{5}=1.40
\end{gathered}
$$

- If we include the two degrees of freedom associated with the vibrational motion (one for the kinetic energy and one for the potential energy), we have

$$
\begin{gathered}
E_{\text {int }}=7 N\left(\frac{1}{2} k_{B} T\right)=\frac{7}{2} N k_{B} T=\frac{7}{2} n R T \\
C_{V}=\frac{1}{n} \frac{d E_{\text {int }}}{d T}=\frac{7}{2} R=29.1 \mathrm{~J} / \mathrm{mol} \cdot K \\
C_{P}=C_{V}+R=\frac{9}{2} R \\
\gamma=\frac{C_{p}}{C_{V}}=\frac{9}{\frac{2}{2} R}=\frac{9}{7}=1.29
\end{gathered}
$$

## Molar Specific Heat of Hydrogen vs. Temperature

Serway/Jewett; Principles of Physics, 3/e
Figure 17.16


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## A Hint of Quantum Physics

- Our classical model can't explain the temperature dependence of $C_{V}$ shown on the previous slide
- To understand this behavior, we need to apply quantum theory in which energy is quantized
- The energy level diagram for the rotational and vibrational states of a diatomic molecule is shown below
- At low temperatures, all the molecules are in the ground state for rotation and vibration, and only the translational motion contributes to the molecules' average energy
- At higher temperatures, molecules are excited into rotational excited states and rotational motion contributes
- At still higher temperatures, vibrational excited states are populated and vibrational motion contributes



## Example 1

A bubble of 5.00 mole of helium is submerged at a certain depth in liquid water when the water (and thus the helium) undergoes a temperature increase $\Delta T$ 0f $20.0^{\circ} \mathrm{C}$ at constant pressure. As a result, the bubble expands. The helium is monatomic and ideal. (a) How much energy is added as heat to the helium during the temperature increase and expansion? (b) What is the change in the internal energy of the helium? (c) How much work is done on the helium?

## Example 2

One mole of a diatomic ideal gas at a temperature of 310 K expands adiabatically from an initial volume of 12.0 L to a final volume of 19.0 L . What is the final temperature of the gas?

Homework Set 6 - Due Wed. Jan. 21

- Read Sections 17.7-17.9
- Do Problems 17.32, 17.35, 17.36, 17.40 \& 17.42

